# Exploring the QCD phase diagram with heavy-ion collisions

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### Outline

- introduction
- theory vs. experiment
- cumulants, factorial cumulants, factorial moments
- expectations
- measurements and interpretation
- summary

#### Quark-gluon plasma as a new state of matter



 $T \sim 300$  K



 $T \sim 1500$  K

Gas of gold atoms:  $T \sim 3000$  K

To melt protons:  $T \sim 10^{12}$  K

Center of the Sun:  $T \sim 10^7$  K



#### Lattice QCD calculations (LQCD)



A smooth crossover

 $170 \text{ MeV} \approx 2 * 10^{12} \text{ K}$ 

Y.Aoki, G.Endrodi, Z.Fodor, S.D.Katz, K.K.Szabo, Nature 443 (2006) 675 A.Bazavov, T.Bhattacharya, M.Cheng et al., Phys. Rev. D80, 014504 (2009) How to create and measure such temperature?





LHC – Large Hadron Collider
 RHIC – Relativistic Heavy Ion Collider, Nowy Jork, USA
 GSI Helmholtz Centre for Heavy Ion Research, Darmstadt, Niemcy

## The collision creates a "quark-gluon plasma"



#### The QCD phase diagram



#### Figure from V.Koch

#### Expectations



Figure from Phys. Rept. 853 (2020) (AB, S.Esumi, V.Koch, J.Liao, M.Stephanov, N.Xu)

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#### Critical point?



A. Pandav, D. Mallick, B. Mohanty, 2203.07817 M. Stephanov, hep-lat/0701002 How to approach this problem? Consider water vapour transition



#### right at the phase transition



number of  $H_2O$  molecules

#### so we measure multiplicity distributions

In QCD we use, e.g., net-baryon, net-charge, net-strangeness

A finite volume van der Waals model



AB, V. Koch, D. Oliinychenko, J. Steinheimer, PRC 98 (2018) 054901

Theory vs. experiment

#### Theory

Coordinate space Fixed volume Long-lived Conserved charges

#### Experiment

Momentum space Expanding and fluctuating volume Extremely short-lived Non-conserved numbers

Lots of detector problems + various corrections (e.g. volume fluctuation)

#### So we measure multiplicity distributions



Au+Au Collisions

- 0.4 < p<sub>T</sub> < 2.0 (GeV/c) |y| < 0.5
- 0-5%

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• 70-80%
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raw distributions (not corrected)
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STAR Collaboration, PRC 104 (2021) 2, 024902

For baryons absolutely minimal goal is to see any deviations from Poisson (Skellam) distribution.

It is difficult to see something in multiplicity distributions. If there is any signal, it is likely very tiny.

We usually characterize P(N) by:

- cumulants  $K_n$ 

- factorial cumualnts,  $C_n$  (or  $\hat{C}_n$ )
- factorial moments  $F_n$  (mean number of pairs, triplets, etc.)

Warning. STAR uses opposite notation  $\kappa_n \leftrightarrow C_n$ 

On the experimental side we need to measure various fluctuation observables and hope to see some nontrivial energy or/and system-size dependence

> see, e.g., Stephanov, Rajagopal, Shuryak, PRL (1998) Stephanov, PRL (2009) Skokov, Friman, Redlich, PRC (2011)

There are many results:

ALICE, STAR, HADES

Cumulants, factorial cumulamts

Proton  $v_1$  (STAR)

HBT radii (STAR) R.A. Lacey, PRL 114 (2015) 142301 NA61/SHINE

Intermittency, cumulamnts

Scaled variance

Strongly intensive variables

#### Poisson distribution (no correlations)



 $P(n) = \text{Poisson if } N \to \infty, \ p \to 0, \ Np = \langle n \rangle$ 

cumulants  $\kappa_i = \langle n \rangle$ 

factorial cumulants  $C_i = 0$ 

factorial moments  $F_i = \langle n \rangle^i$ 

Factorial cumulants – example



## m particle cluster



$$C_2 \neq 0$$
$$C_k = 0, k > 2$$

 $C_{2,3,...,m} \neq 0$  $C_k = 0, k > m$ 

factorial  
cumulants 
$$C_k = \frac{d^k}{dz^k} \ln\left(\sum_n P(n)z^n\right)|_{z=1}$$

Poisson

Two-particle correlation function

$$\rho_2(y_1, y_2) = \rho(y_1)\rho(y_2) + C_2(y_1, y_2)$$

Integrating both sides over some bin in rapidity

$$\langle n(n-1)\rangle = \langle n\rangle^2 + C_2$$

$$C_2 = \int C_2(y_1, y_2) dy_1 dy_2$$

Same with multiparticle correlations.

Factorial cumulants are integrated multiparticle correlation functions

Factorial cumulants vs cumulants

factorial cumulant

$$C_i = \frac{d^i}{dz^i} \ln\left(\sum_n P(n) z^n\right)|_{z=1}$$

cumulant 
$$\kappa_i = \frac{d^i}{dt^i} \ln\left(\sum_n P(n)e^{tn}\right)|_{t=0}$$

Poisson  $C_i = 0, \ \kappa_i = \langle n \rangle$  cumulants naturally appear in statistical physics

$$\ln(Z) = \ln\left(\sum_{i} e^{-\beta(E_i - \mu N_i)}\right)$$

Cumulants (one species of particles)

$$\kappa_{2} = \langle N \rangle + C_{2}$$
  

$$\kappa_{3} = \langle N \rangle + 3C_{2} + C_{3}$$
  

$$\kappa_{4} = \langle N \rangle + 7C_{2} + 6C_{3} + C_{4}$$

Cumulants mix integrated correlation functions of different orders

They might be dominated by  $\langle N \rangle$ .

See, e.g., B. Ling, M. Stephanov, PRC 93 (2016) 034915 AB, V.Koch, N.Strodthoff , PRC 95 (2017) 054906 "Cumulant ratios do not depend on volume"

but depend on volume fluctuation

It is true if a correlation length is much smaller than the system size



Here this condition is satisfied

#### momentum rapidity space



Correlation length is usually larger than one unit of rapidity.

Cumulant ratios are expected to depend on acceptance in rapidity

#### Short-range correlations

 $C_i \sim \langle N \rangle \sim \Delta y$  $\kappa_i \sim \langle N \rangle \sim \Delta y$ 

Long-range correlations (expected in rapidity)

$$C_i \sim \langle N \rangle^i \sim (\Delta y)^i$$

 $\kappa_i$  is complicated, for example

$$\kappa_{4} = \langle N \rangle + (\sim \langle N \rangle^{2}) + (\sim \langle N \rangle^{3}) + (\sim \langle N \rangle^{4})$$
  

$$\kappa_{2} = \langle N \rangle + (\sim \langle N \rangle^{2})$$
polynomial in  $\Delta y$ 

Cumulant ratios may strongly depend on acceptance in rapidity and in transverse momentum

Comparison with models which do not have experimental acceptance is questionable

Comparison with lattice QCD calculations is very tricky

With long-range rapidity correlations the cleanest observable is  $\frac{C_i}{I M N_i}$ 



HRG with attractive and repulsive Van der Waals interactions between (anti)baryons



V. Vovchenko, M.I. Gorenstein, H. Stoecker, PRL 118 (2017) 182301



Density plot of the quartic cumulant obtained by mapping the Ising model into QCD. Freezeout line is for demonstration only.



Normalized quartic cumulant of proton multiplicity



D. Mroczek, A.R. Nava Acuna, J. Noronha-Hostler, P. Parotto, C. Ratti, M.A. Stephanov, PRC 103 (2021) 3, 034901

We find that, while the peak remains a solid feature, the presence of the critical point does not necessarily cause a dip in  $\chi_4^B$  on the freezeout line below the transition temperature.

#### **STAR Collaboration**

2112.0024 PRC 104 (2021) 2, 024902



Visible four-proton correlations at 7.7 GeV (large errors)

A hint of non-monotonic dependence

STAR data vs. hydrodynamics with baryon conservation and excluded volume



Baryon conservation for  $\sqrt{s} > 20$  GeV

V.Vovchenko, V.Koch, C.Shen, PRC 105, 014904 (2022) P.Braun-Munzinger, B.Friman, K.Redlich, A.Rustamov, J.Stachel, NPA 1008 (2021) 122141 AB, V.Koch, V.Skokov, PRC 87 (2013) 1, 014901

#### **ALICE Collaboration**

#### PLB 807 (2020) 135564



**Global** (not local) baryon conservation! Something to understand. It would be good to measure proton, antiproton and mixed proton-antiproton factorial cumulants M.Barej, AB, PRC 102 (2020) 6, 064908

See O.Savchuk, V.Vovchenko, V.Koch, J.Steinheimer, H.Stoecker, PLB 827 (2022) 136983 Local conservation and  $B\overline{B}$  annihilation **HADES** Collaboration

PRC 102 (2020) 2, 024914



### NA61/SHINE Collaboration



#### NA61/SHINE Collaboration



No critical signal. Consistent with p+p.

#### Conclusions

Interesting and important physics but so far no success

Clear signal of (global?) baryon conservation

Interesting STAR point at 7.7 GeV. Four-proton correlations (physics?)

We definitely need better statistics

Hopefully some progress will come from lattice QCD

